

# Determinantal Quantum Monte Carlo solver for Cluster Perturbation Theory

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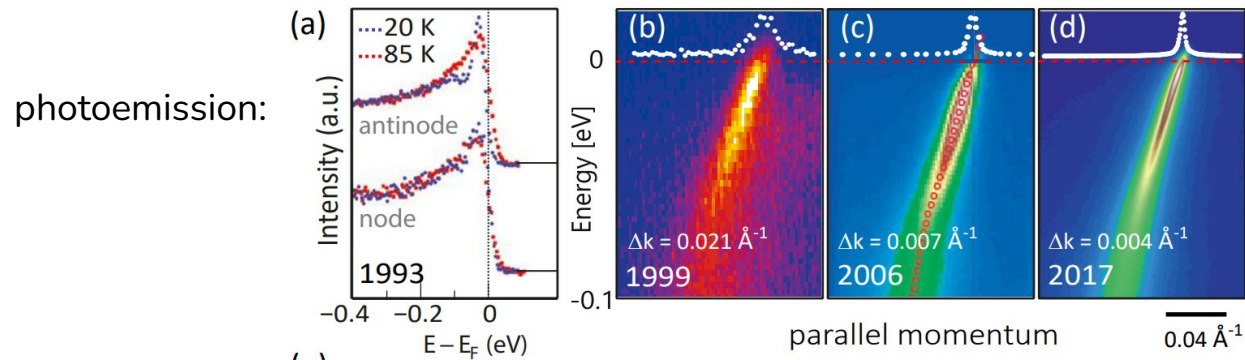


$A(k, \omega)$

$$A^{<}(k, \omega) = \sum_n |\langle n | c_k | G \rangle|^2 \delta(\omega + (E_n - E_G))$$

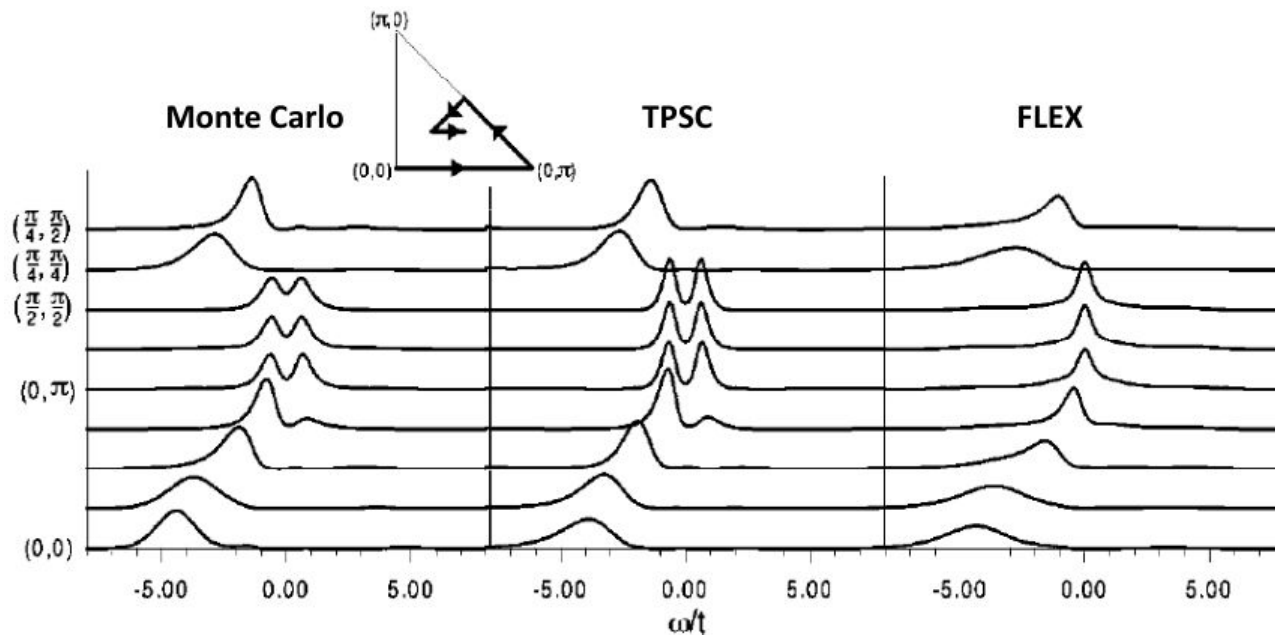
$$A^{>}(k, \omega) = \sum_n |\langle n | c_k^\dagger | G \rangle|^2 \delta(\omega - (E_n - E_G))$$

$$A^{<}(k, \omega) + A^{>}(k, \omega) = A(k, \omega)$$



# How to calculate?

- Semi-analytical/perturbative methods (e.g. FLEX)
  - Low computational cost, high momentum resolution, low temperatures.
  - Approximate. Questionable validity for intermediate/strongly interacting systems.
- Finite cluster methods (e.g. DQMC, ED)
  - Exact.
  - High computational cost. Limited cluster size/momentum resolution, limited temperatures.

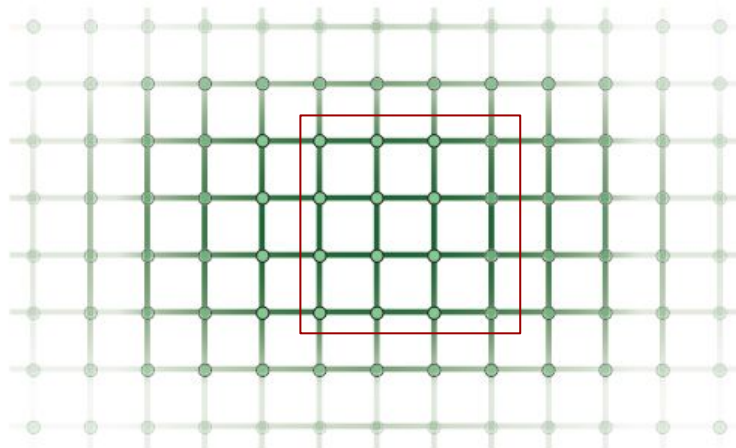


$U/t=4, \langle n \rangle=1, T/t=0.2$

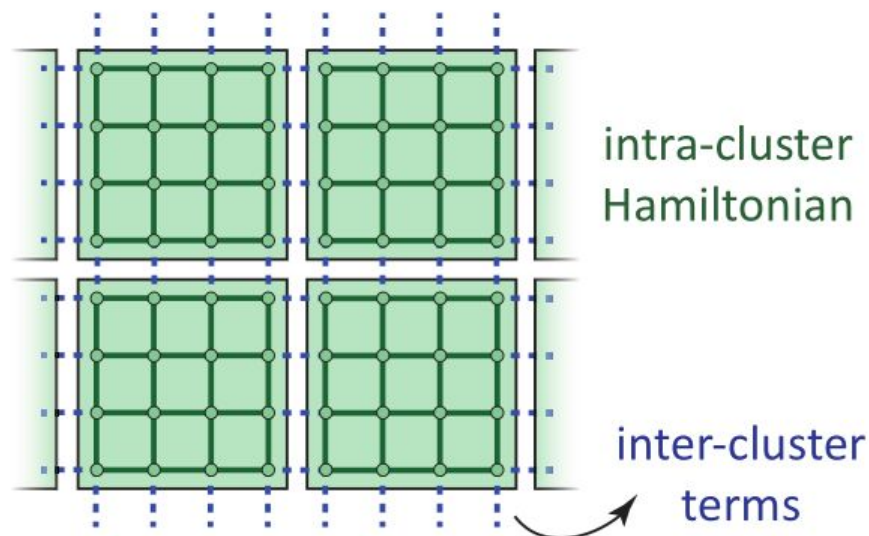
Moukouri et al, PRB 2000

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  - High computational cost. Limited cluster size/momentum resolution, limited temperatures.
- Embedding methods (e.g. CDMFT, DCA, CPT)
  - “Exact at short-range, approximate at long-range”



# Cluster perturbation theory (CPT)



$$\mathcal{H} = \sum_{\mathcal{C}} \mathcal{H}^{\mathcal{C}} + \sum_{ij\sigma} h_{ij}^b c_{i\sigma}^{\dagger} c_{j\sigma}$$

- [12] S. Pairault, D. Sénéchal, and A.-M. Tremblay, Phys. Rev. Lett. **80**, 5389 (1998).
- [13] D. Sénéchal, D. Perez, and M. Pioro-Ladrière, Phys. Rev. Lett. **84**, 522 (2000).
- [14] D. Sénéchal, D. Perez, and D. Plouffe, Phys. Rev. B **66**, 075129 (2002).
- [15] D. Sénéchal, in *Strongly Correlated Systems* (Springer, 2012) pp. 237–270.

# Cluster perturbation theory (CPT)

$$\mathcal{H} = \sum_c \mathcal{H}^c + \sum_{ij\sigma} h_{ij}^b c_{i\sigma}^\dagger c_{j\sigma}$$

$$\mathbf{G}(z) = \frac{\mathbf{G}^c(z)}{\mathbf{I} - \mathbf{h}^b \mathbf{G}^c(z)}$$

- Exact at  $U=0$ .

$$\mathbf{G}^c(z) = \frac{1}{z - \mathbf{h}^c}$$

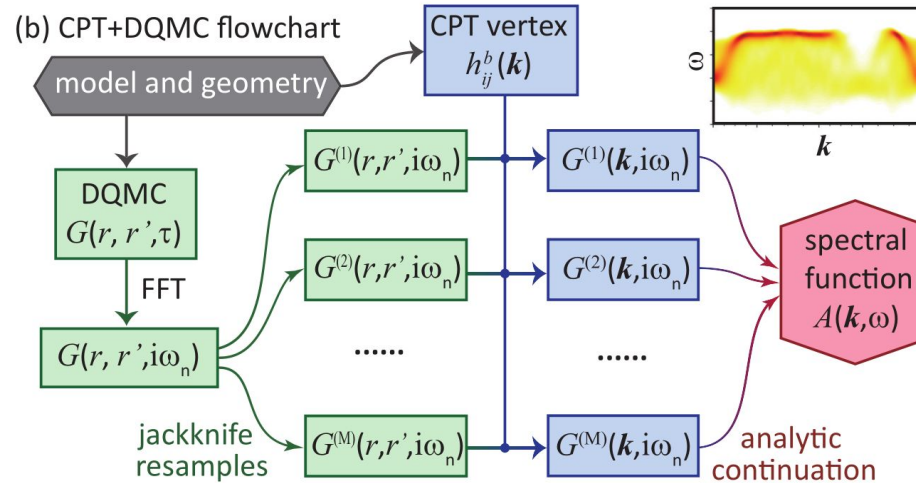
$$\mathbf{G}(z) = (z - \mathbf{h}^c - \mathbf{h}^b)^{-1} = \frac{1}{z - \mathbf{h}}$$

- Exact at  $t=0$ .

# Solving $H^c$ to obtain $G^c$

- Exact diagonalization (ED):
  - Limited to  $\sim 16$  site clusters.
  - Zero temperature. Finite temperature limited to fewer sites.
- Time-dependent DMRG (Yang, Feiguin, PRB 2016)
  - Restricted to narrow ladder geometries
- DQMC:
  - Larger system sizes ( $\sim 100$ s) possible.
  - Finite temperature
  - Possible to handle problems difficult for ED (e.g. Holstein model)

# DQMC solver for CPT (CPT+DQMC)



1. DQMC:  $G_{rr'}^c(\tau) = - \left\langle c_{r\sigma}(\tau) c_{r'\sigma}^\dagger \right\rangle$
2. Fourier transform:  $G_{rr'}^c(i\omega_n) = \int_0^\beta d\tau e^{i\omega_n \tau} G_{rr'}^c(\tau)$
3. Jackknife resample.
4. CPT.  $G(k, z) = \frac{1}{N^c} \sum_{r, r'} e^{-ik(r-r')} \left[ \frac{\mathbf{G}^c(z)}{\mathbf{I} - \tilde{\mathbf{h}}^b \mathbf{G}^c(z)} \right]_{r, r'}$
5. MaxEnt  $G(k, i\omega_n) = \int d\omega \frac{A(k, \omega)}{i\omega_n - \omega}$

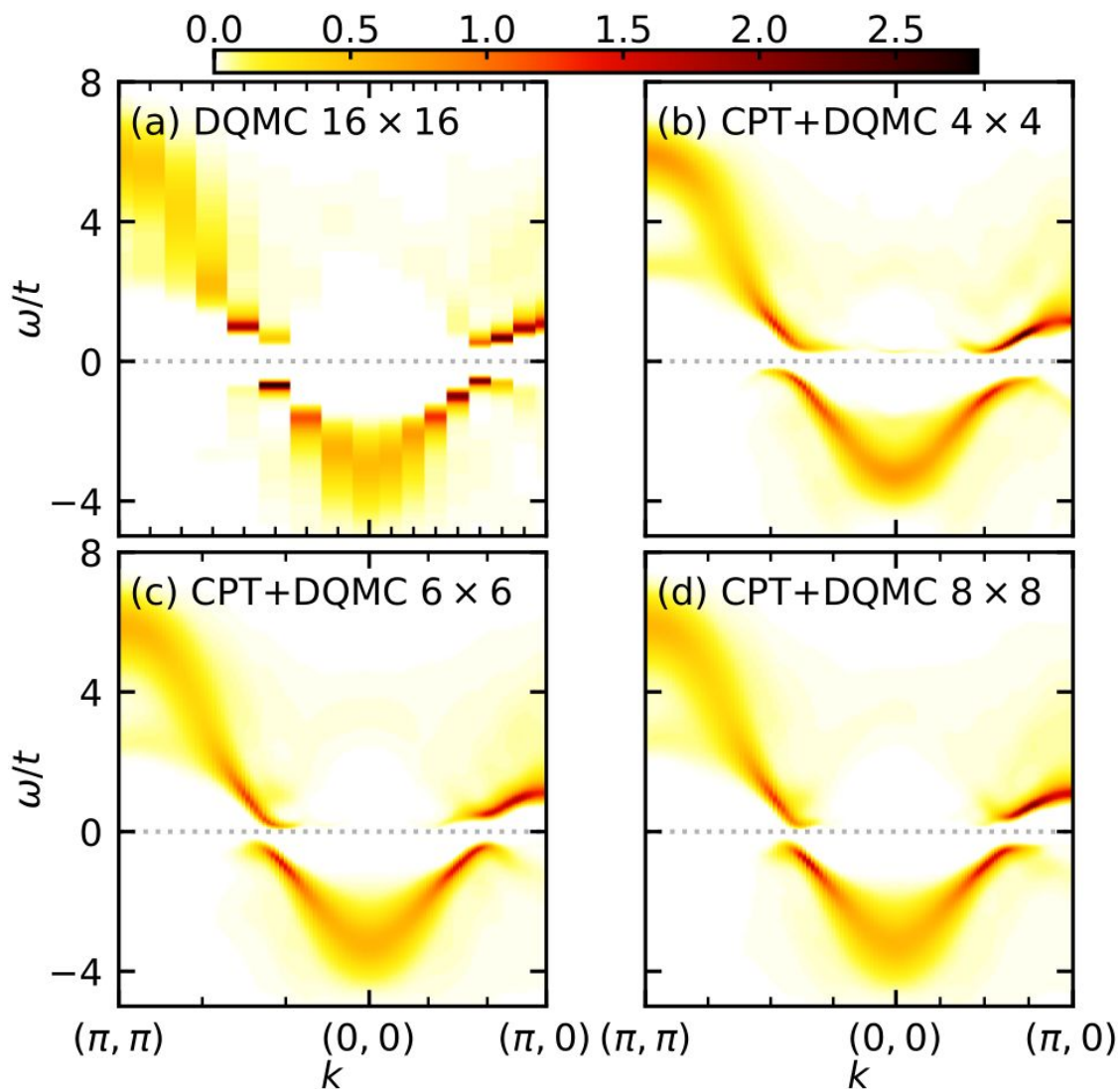


# Benchmarking: attractive Hubbard model, SC phase

- Sign-free DQMC simulations.
- s-wave SC with KT transition at  $T_c/t \approx 0.14$  for  $U/t = -4$ ,  $\langle n \rangle = 0.6$

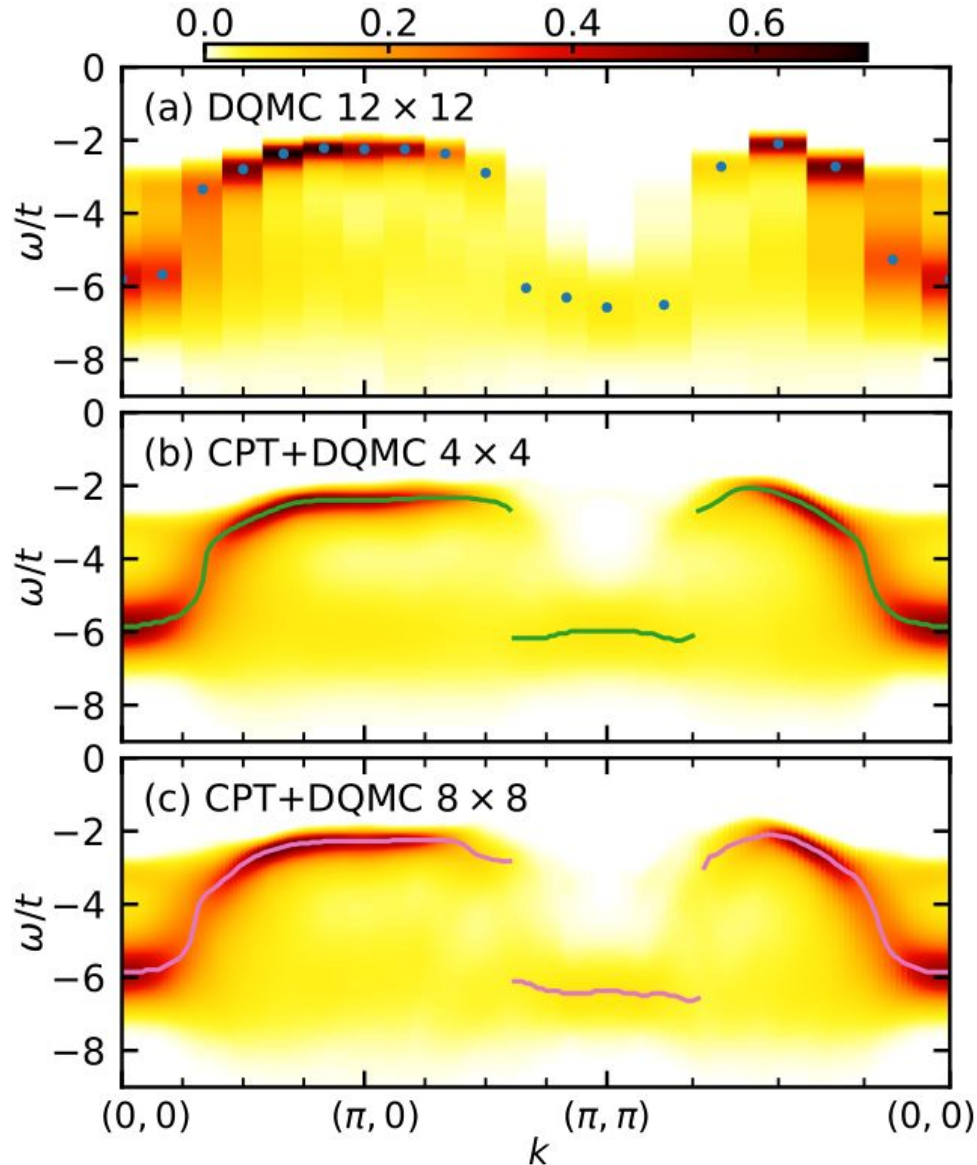
# Benchmarking: attractive Hubbard model, SC phase

$U/t = -4$   
 $\langle n \rangle \approx 0.63$   
 $T/t = 1/12$



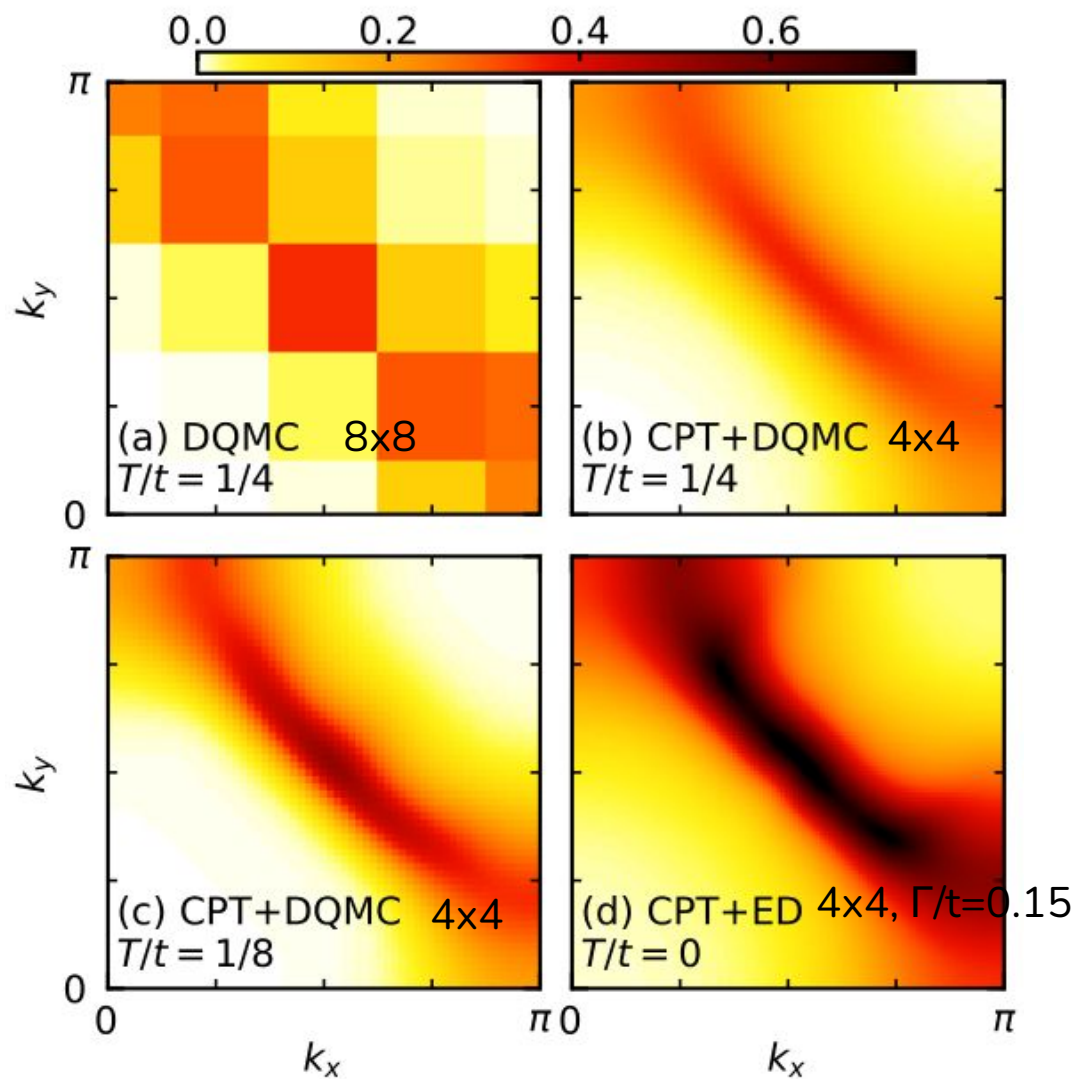
# Benchmarking: half-filled repulsive Hubbard model

$U/t = 8$   
 $\langle n \rangle = 1$   
 $T/t = 1/16$



# Doped repulsive Hubbard model: $A(k, \omega = 0)$

$U/t = 8$   
 $\langle n \rangle \approx 0.94$



# Summary

- Advantages over DQMC:
  - Continuous momentum resolution
  - Smaller clusters
    - Reduced sign problem
    - Lower computational cost
- Advantages over CPT+ED:
  - Larger system sizes/reduced finite-size effects
  - Temperature dependence
  - Handles multi-orbital models, el-ph models, etc more easily

Next steps:

- Pseudogap
- Apply better methods of analytic continuation
- Apply to more models